The role of tropical analysis uncertainties in global predictability





Motivation: 12-hour forecast uncertainties



Flow dependent uncertainties

150 hPa, tropics



370 hPa, midlatitudes

- Spread in -12hour zonal wind ensemble (in m/s) along the latitude circle
- 3-month long experiment with a perfect model and 12-hour cycle EnKF data assimilation

MWR, 2016

Flow dependency of the simulated forecast errors in EDA



3-h fc errors in the zonal wind, derived from the ECMWF ensemble (cy32r3) during 1 month (July 2007)

QJRMS, 2013

Growth of forecast uncertainties in ensemble prediction





- ECMWFENS
 - Two weeks of data in May 2015
- Ensemble spread in the zonal wind

lev45, ~150 hPa lev55, ~290 hPa



Zonally-averaged forecast-error statistics



ECMWF ensemble prediction system: two weeks of data in May 2015 Ensemble spread in zonal wind (m/s)

Tropical analysis uncertainties: summary

Uncertainties in global analyses and short-range forecasts are largest in the tropics

- + A lack of observations, especially wind observations
- + Complex moist dynamics
- Data assimilation methodologies focused on the extratropics

Tropics and the global observing system



Impact of tropical analysis uncertainties on the midlatitude forecasts

Heating perturbation (up to 0.5 K/day) over Indian ocean and Maritime continent



- wind and geopotential height perturbations at 200 hPa level
- average of 30 simulations started on 1 January
- SPEEDY general circulation model, T₃₀

PhD thesis research by Katarina Kosovelj

Decomposition of tropical heating perturbations

1-day average response to the large-scale heating perturbation



Decomposition of the global response to the tropical heating perturbations

14-day average response at 200 hPa to the large-scale heating perturbation Total Balanced





Unbalanced

Kelvin wave

Balanced n=1 Rossby



Decomposition of forecast errors (12-hour ensemble spread)



- Spread in 12-hr zonal wind forecast, model level close to 250 hPa
- The WIG spread is greater in the mid-latitudes in relation to waves developing on the mean westerly flow.
- The EIG component is larger than the WIG spread in the tropics

MWR, 2016

Distribution of tropical forecast-error variance among equatorial modes



Impact of the equatorial wave constraint on analysis increments

Single h observations at the equator







Kelvin wave coupling is decisive for the structure of analysis increments near the equator

Other equatorial inertio-gravity waves reduce the meridional correlation scale, as well as effect the mass-wind coupling

QJRMS, 2005

Impact of the equatorial wave constraint on analysis increments

Single westerly wind obs at the EQ

Rossby waves



Tropical moist dynamics and 4D-Var

Single moisture observation at the equator, 12-h 4D-Var



Simple (modelling) is beautiful

Coupling between winds, moisture and aerosols in 4D-Var PhD thesis research by Ziga Zaplotnik

Tropical moist dynamics and 4D-Var

Single wind observation at the equator, 12-h 4D-Var



Representation of the global forecast-error variances using the Hough functions

$$\mathbf{B} \approx \frac{1}{N_{\text{ens}} - 1} \sum_{i=1}^{N_{\text{ens}}} \Delta \mathbf{x}_i^B \cdot (\Delta \mathbf{x}_i^B)^{\text{T}}$$

$$\mathbf{L} = \mathbf{D}\boldsymbol{\Theta}_{\boldsymbol{\varphi}}\mathbf{F}_{\boldsymbol{\lambda}}\mathbf{G}_{m}$$

 $\chi = \mathbf{L}\Delta \mathbf{x}^{B}$

Estimate of the bkg error from the ensemble

- G_y projection on the vertical structure Θ – projection on the meridionally part of Hough harmonics
- D- spectral variance density normalization
- F Fourier transform in the zonal direction

Entropy reduction

Forecast-error variance reduction

$$S_{\chi} = \frac{1}{2} \sum_{\nu=1}^{N} \ln \mathcal{I}_{\nu}^{-1}$$

M. Fisher, 2003

 $\mathcal{I}_{\nu} = \mathcal{I}_{n}^{k}(m) = \frac{[\gamma_{\nu}^{A}]^{2}}{[\gamma^{B}]^{2}}$

MWR, 2016

Analysis and forecast uncertainties in Observing System Simulation Experiment with a perfect model

Data Assimilation Research Testbed (DART), by Jeff Anderson and collaborators, *http://www.image.ucar.edu/DAReS/DART/*

Spectral T85 Community Atmosphere Model, CAM 4 physics

Long spin-up (from 1 Jan 2008) with the observed SST to reproduce nature run ('truth')

Preparation of the observations from the nature run

Preparation of the homogeneous observing network (Δ ~920 km)

Assimilation cycle during three months (Aug-Oct) in 2008

No inflation

Decomposition of forecast errors (12hour ensemble spread)



- Spread in 12-hr zonal wind forecast, model level close to 250 hPa
- The WIG spread is greater in the mid-latitudes in relation to waves developing on the mean westerly flow.
- The EIG component is larger than the WIG spread in the tropics

MWR, 2016

Growth of the global forecast uncertainties

Zonally-averaged zonal wind spread in ensemble of 12-hr forecasts



Perfect-model experiment

Growth of the global forecast uncertainties

Zonally-averaged zonal wind spread in ensemble of 12-hr forecasts



Maximum of analysis and short-term forecast uncertainties in the upper tropical troposphere is not due to model error

Spread of the analysis ensemble



m/s

Reduction of the forecast uncertainties by the assimilation



- Prior (12-hr forecast) posterior (analysis) spread, normalized by the prior spread
- Spread reduction is greater in the mid-latitudes
- Spread is poorly reduced in the tropics

Short-range global forecast errors in the perfect-model EnKF framework



Spread of 12-hr forecast ensemble 3-month average

(Prior – posterior)/prior ensemble spread (x,y,z) points averaged in time and zonally

%

0.18 0.16

0.14

0.12 0.1

0.08

0.06

0.04

0.02

80

60

Data assimilation efficiency: variance reduction



Efficiency = variance reduction as a function of zonal wavenumber

The assimilation is most efficient in synoptic scales, for both balanced and IG motions but much more efficient for balanced.

Covariance localization radius was 0.2 (around 1300 km at Eq).

MWR, 2016

Impact of the covariance localization radius



Scale-dependent growth of the global forecast uncertainties towards saturations

Forecast was started on 1 Oct 2008 in a perfect model EnKF OSSE





Ensemble spread in different zonal wavenumbers is normalized by its initial value

Log(E(k,t)/E(k,o)]

Ensemble spread in each zonal wavenumber is normalized by its value at 50-day forecast range.

Scale-dependent growth of the global forecast uncertainties towards saturations in ENS



Summary and outlook

- Tropics are characterized by largest analysis uncertainties and largest growth of forecast uncertainties during the first 24-36 hours of the forecast
- The uncertainties are on average larger on the large scales. Maximum of uncertainties is in the tropical upper troposphere.
- + Uncertainties are flow dependent. Uncertainties in wind and geo. height fields in the tropics are balanced about 50%.
- In an OSSE with a perfect mode and EnKF, the covariance localization radius is important in the tropics.
- Flow-dependent ensemble -> fc-error variance spectrum of the day -> weights for the mass-wind constraint in the bkg-error term for various IG modes and Rossby modes of the day

Additional slides



Tropical data assimilation system including Rossby and IG wave constraints

 Application of parabolic cylinder functions as the basis functions for the representation of the background-error covariances

$$J(\chi) = J_{b} + J_{o} = \frac{1}{2}\chi^{\mathrm{T}}\chi + \frac{1}{2}\sum_{n=1}^{K} (\mathbf{y}_{n} - \mathbf{H}(\mathbf{x}^{b} + \mathbf{L}^{-1}\chi_{n}))^{\mathrm{T}}\mathbf{R}^{-1}(\mathbf{y}_{n} - \mathbf{H}(\mathbf{x}^{b} + \mathbf{L}^{-1}\chi_{n}))$$

$$\chi = L \delta \mathbf{x}$$

$$L = DP_y F_x F^{-1}$$

 P_y – projection operator on the meridionally dependent part of equatorial eigenmodes D – spectral variance density normalization F – Fourier transform operator



History of Hough functions in data assimilation (1)

- + Flattery, 1970s: NCEP OI based on the Hough functions
- + D. Parrish, mid 1980s: computed correlations for single point in the tropics including the impact of KW and MRG waves

Single height observations at EQ







Parrish, 1988, AMS proceedings



(h,h), Rossby+MRG

(h,h), Rossby+MRG+KW, k=1-3 (h,u), Rossby+MRG+KW, k=1-3

History of Hough functions in data assimilation (2)

+ ECMWF, early 1990s: first formulation of 3D-Var used Hough functions

$$J_{b} = \frac{1}{2} c_{R} \chi_{R}^{t} \Lambda_{R}^{-1} \chi_{R} + \frac{1}{2} c_{G} \chi_{G}^{t} \Lambda_{G}^{-1} \chi_{G} + \frac{1}{2} \chi_{U}^{t} \Lambda_{U}^{-1} \chi_{U}$$

$$c_{G} \text{ is set to } \frac{1}{2\epsilon} \text{ and } c_{R} \text{ to } \frac{1}{2(1-\epsilon)}$$
Heckley et al., 1993, ECMWF proceedings



wind obs at the EQ

Single easterly wind obs at the EQ

Single southerly wind obs at the EQ

History of Hough functions in data assimilation (3)

+ ECMWF, early 1990s: first formulation of 3D-Var used Hough functions



Single westerly wind obs at the EQ at 500 hPa

Potential impact of ADM-Aeolus in the tropics: Rossby wave example



pears index

Žagar et al., MWR 2008

Growth of the global forecast errors in the perfect-model

Perfect-model experiment



Zonally-averaged zonal wind spread in ensemble of 12-hr forecasts



(Forecast– analysis)/forecast ensemble spread in each (x,y,z), averaged in time and zonally

Data assimilation efficiency: variance reduction



The assimilation is most efficient in synoptic scales, for both balanced and IG motions

MWR, 2016

Scale-dependency of the 12-hr forecast error variances in EnKF with a perfect model



Distribution of the variance in analysis ensmeble looks very similar. As expected, largest variance is in synoptic scales and balanced modes (mid-latitudes) and in the large-scale Kelvin wave

Žagar et al., 2016, MWR

Flow dependent growth of forecast uncertainties in ENS



Scale-dependent growth of forecast uncertainties: classical approach

Partition of ECMWF ENS spread at 200 hPa level into rotational and divergent parts



Zonally-averaged ensemble spread in EDA





Žagar et al., 2013, QJRMS

Short term growth of simulated forecast errors in EDA in relation to flow



In the tropics, the short-range growth is largest in the Kelvin mode The growth in WIG modes is accompanying the balanced variance growth in the midlatitudes

Scale and flow dependent representation of the ensemble reliability

Dec 2014, Operational ECMWF ENS data

A lack of variability is initially seen in subsynoptic balanced scales, and lateron in tropical IG modes, primarily the Kelvin mode



Žagar et al., 2015, JAS

Growth of the ensemble spread w.r.t. initial spread as a function of zonal scale



Initially, spread growth is largest in the smallest scales and the synoptic scales of the IG modes (tropics)

Žagar et al., 2015, JAS

Growth of the IG spread w.r.t. initial spread



Zadar et al., 2015, JAS

Short-range forecast error statistics, EDA



Almost half of the variance in short-term forecast errors is associated with the inertio-gravity modes. EIG dominates over WIG on all scales. Data from July

1D growth of forecast uncertainties

Growth of uncertainties in m=1 vertical mode in ECMWF ENS in May 2015



Žagar, 2017, Tellus A

Scale-dependent limits of the growth of spread in ENS





Dalcher and Kalnay, 1987

Growth of error variance for Z500 in the ECMWF model in early 1980s. The smaller the scale, the shorter the predictability limit

Scale-dependent growth of forecast uncertainties



Žagar, 2017, Tellus A

Growth of the spread w.r.t. initial spread as a function of the zonal scale



(industrial space at the equation of the

Scale-dependent growth of forecast uncertainties

Based on model-level data from the operational ECMWF ENS in May 2015



- Global analysis and forecast uncertainties in the first vertical mode as a function of the zonal wavenumber and meridional mode
- Only spread associated with Rossby modes (balanced dynamics)

Growth of the 3D integrated uncertainties



3D integrated ensemble spread shown as log E(t) on various scales during 15 days of operational ECMWF ENS data in May 2015

Chaotic nature of atmospheric system is evidenced in the exponential growth of errors

Žagar, 2017, Tellus A

Fitting the growth of the integrated ensemble spread



Normalize data by initial spread

A new function fit to data that provides analytical estimate of the asymptotic curves Scale-dependent growth of the global forecast errors towards saturations

Simulated forecast errors in different zonal wavenumbers normalized by their asymptotic values at 60-day forecast range



Žagar et al., 2017, Tellus A

Scale-dependent growth of the global forecast errors towards saturations

Simulated forecast errors in different zonal wavenumbers normalized by their asymptotic values at 60-day forecast range



Žagar et al., 2017, Tellus A

Tropics: impact on the midlatitudes

Impact of tropical heating perturbations on midlatitudes



- meridional wind perturbations along 50N at 200 hPa level,
- average of 30 simulations started on 1 January with heating perturbation in the Indian ocean/Maritime continent
- SPEEDY general circulation model

ongoing PhD research by Katarina Kosovelj

Uncertainty partition into balanced and unbalanced components



Uncertainty partition into balanced and unbalanced components



Scale-dependent representation of analysis and forecast uncertainties



Balanced part of circulation is associated with the Rossby (quasi-geostrophic) part of eigensolutions to the linearized primitive equations. The unbalanced part projects onto the inertio-gravity eigensolutions that propagate eastward (EIG modes) or westward (WIG modes).



MODES, http://meteo.fmf.uni-lj.si/MODES

MODES People -28.11.2016 00UTC +000h IG = 8% BAL= 92% 10 10³ Total energy (J/kg) Modal view of atmospheric circulation 10⁰ MODES focuses on the representation of the inertio-gravity circulation in numerical weather p ensemble prediction systems and climate simulations. The project methodology relies on the c circulation in terms of 3D orthogonal normal-mode functions. It allows quantification of the ro TOTAL atmospheric varibility across the whole spectrum of resolved spatial and temporal scales. 10 IG BAL

http://meteo.fmf.uni-lj.si/MODES

10

2

3

25

10

15

Zonal wavenumber k

50

100

180

MORE ABOUT MODES

Expansion of discrete data: vertical projection

An input data vector \mathbf{X} is defined on the horizontal regular Gaussian grid and vertical sigma levels at time t.

$$\mathbf{X}(l,j,S) = (u,v,h)^T$$

Projection of a single data point on j-th sigma level is performed on the precomputed vertical structure functions **G**, the horizontal Hough vector functions in the meridional direction and waves in the longitudinal direction:

$$\mathbf{X}(l,j,S_j) = \mathop{\bigotimes}\limits_{m=1}^{M} \mathbf{S}_m \mathbf{X}_m(l,j) \times G_m(j)$$
(1)

The vector \mathbf{X}_m is obtained by the reverse transform of (1):

$$\mathbf{X}_{m}(\boldsymbol{1},\boldsymbol{j}) = \mathbf{S}_{m}^{-1} \overset{J}{\underset{j=1}{\overset{J}{\overset{j}{a}}}} \left(\boldsymbol{u},\boldsymbol{v},\boldsymbol{h}\right)_{j}^{T} \boldsymbol{G}_{m}(\boldsymbol{j})$$
(2)

Two kinds of Hough harmonic solutions for the horizontal wave motions

Frequencies of spherical normal modes for different equivalent depths



Meridional structure of Hough functions

HSFs are pre-computed for a given number of vertical modes, M For every m=1,...,M, i.e. for every D_m

Meridional structure for Hough functions is computed for a range of the zonal wavenumbers K,

k=-K,...,o,...,K

and a range of meridional modes for the balanced, N_{ROSSBY} , a range of EIG, N_{EIG} , and a range of WIG, N_{WIG} , modes.

 $R=N_{ROSSBY} + N_{EIG} + N_{WIG}$



Žagar et al., 2015, GMD

Expansion of discrete data: horizontal projection

The horizontal coefficient vector \mathbf{X}_m for a given vertical mode is projected onto the Hough harmonics $\mathbf{H}^n_k(\lambda, \varphi, m)$ as

$$\mathbf{X}_{m}(\mathbf{1},\mathbf{j}) = \mathop{\mathbb{a}}\limits_{n=1}^{R} \mathop{\mathbb{a}}\limits_{k=-K}^{K} C_{n}^{k}(m) \mathbf{H}_{n}^{k}(\mathbf{1},\mathbf{j},m)$$
(3)

The subscript *n* indicates all meridional modes including rotational (ROT), and eastward and westward propagating inertio-gravity (EIG and WIG, respectively) modes

The scalar complex coefficients χ are obtained as

$$\chi_n^k(m) = \frac{1}{2\pi} \int_{0}^{2\pi} \int_{-1}^{1} \left(\tilde{u}_m, \tilde{v}_m, \tilde{h}_m \right)^T \left[\mathbf{H}_{n'}^{k'} \right]^* d\mu d\lambda$$
(4)

Here, $\mu = sin(\phi)$.

Expansion of discrete data: energy product

The partition of total energy into the kinetic and available potential energy for every vertical mode is written as :

$$\frac{\P}{\P t} \stackrel{2\rho}{\overset{1}{0}} \stackrel{1}{\overset{0}{0}} \stackrel{a}{\underset{m}{\partial}} \frac{1}{2} \stackrel{a}{\overset{p}{e}} u_{m}^{2} + v_{m}^{2} + \frac{g}{D_{m}} h_{m}^{2} \stackrel{\ddot{0}}{\overset{+}{\Rightarrow}} a^{2} dm d/ = 0$$

Global energy product of the m-th vertical mode defined as

$$I_{m} = \frac{1}{2}gD_{m} \overset{R}{\overset{R}{a}} \overset{K}{\overset{R}{a}} C_{n}^{k}(m) \overset{L}{\overset{K}{b}} C_{n}^{k}(m) \overset{L}{\overset{K}{b}}^{*}$$

is equivalent to

$$I_{m} = \frac{1}{2}gD_{m} \overset{2\rho}{\overset{0}{\underset{0}{\circ}}} \overset{1}{\underset{-1}{\circ}} \left(\tilde{u}_{m}^{2} + \tilde{v}_{m}^{2} + \tilde{h}_{m}^{2} \right) d / d m = \overset{2\rho}{\underset{0}{\circ}} \overset{1}{\underset{-1}{\circ}} \left(K_{m} + P_{m} \right) d / d m$$

Ensemble spread in modal space

If the input fields to the projection are differences between the ensemble members n=1,..,N and the ensemble mean, the total variance in the modal space is defined as

The specific modal variance Σ^2 is defined as

$$\underbrace{\hat{\beta}} S_n^k(m) \underbrace{\hat{\beta}}^2 = \frac{1}{P - 1} \underbrace{\hat{\beta}}_{p=1}^p g D_m \left(C_n^k(m; p) \underbrace{\hat{\beta}} C_n^k(m; p) \underbrace{\hat{\beta}}^* \right)$$

The modal-space variance defined by (5) is equivalent to the total variance in the physical space defined as

 $\underset{i \quad j \quad m}{\overset{a}{\underset{j}}} \overset{a}{\underset{m}} S^{2}(I_{i}, j_{j}, m)$

 $aaa B S_n^k(m)$

k n m

with the specific variance in physical space
$$\mathbf{S}^2$$

$$S^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) = \frac{1}{P-1} \bigotimes_{p=1}^{P} \bigotimes_{e}^{\mathcal{A}} u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + v_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + \frac{g}{D_m} h_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) \stackrel{\text{"orange}}{\div} \bigotimes_{\emptyset}^{\mathcal{O}} u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + v_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + \frac{g}{D_m} u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) \stackrel{\text{"orange}}{\div} \bigotimes_{\emptyset}^{\mathcal{O}} u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) \stackrel{\text{"orange}}{\to} u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) \stackrel{\text{"orange}}{\to} u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) + u_p^2(\boldsymbol{I}_i, \boldsymbol{j}_j, \boldsymbol{m}) \stackrel{\text{"orange}}{\to} u_p^2(\boldsymbol{I}_j, \boldsymbol{J}_j$$

Žagar et al., 2015, JAS

Inertio-gravity circulation of the day

